

Lecture 1

Heisenberg spin chain

A mini-course “Solvable lattice models and Bethe Ansatz”
(Ariel University, spring 2021)

Michael Lashkevich

Landau Institute for Theoretical Physics,
Kharkevich Institute for Information Transmission Problems

Single spin $1/2$

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It satisfies the condition

$$T\sigma_n^\alpha = \sigma_{n-1}^\alpha T. \quad (8)$$

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$$H_{XYZ} = -\frac{1}{2} \sum_{n=1}^N \left(J_x \sigma_n^x \sigma_{n+1}^x + J_y \sigma_n^y \sigma_{n+1}^y + J_z \sigma_n^z \sigma_{n+1}^z \right) \quad (9)$$

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With the cyclic boundary condition the Hamiltonian is translationally invariant. It commutes with the translation operator:

$$[H_{XYZ}, T] = 0. \quad (10)$$

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Each term either does not change the spins or flips one spin up and another spin down. Hence

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Therefore we can split the space of states into the subspaces of spin eigenvectors:

$$\mathcal{H}_N = \bigoplus_{k=0}^N (\mathcal{H}_N)_k, \quad (\mathcal{H}_N)_k = \{ |\psi\rangle \mid S^z |\psi\rangle = (N/2 - k) |\psi\rangle \}. \quad (14)$$

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The Hamiltonian H_{XXZ} acts as an operator on each of these subspaces.

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The space $(\mathcal{H}_N)_2$ is $N(N-1)/2$ -dimensional:

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Generally we have

$$(\mathcal{H}_N)_k = \bigoplus_{1 \leq n_1 < \dots < n_k \leq N} \mathbb{C}|n_1, \dots, n_k\rangle,$$

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The space $(\mathcal{H}_N)_0$ is one-dimensional:

$$(\mathcal{H}_N)_0 = \mathbb{C}|\Omega_+\rangle, \quad |\Omega_+\rangle = |\uparrow\uparrow\dots\uparrow\rangle. \quad (15)$$

This state can be also defined by the condition

$$\sigma_n^+ |\Omega_+\rangle = 0 \quad \forall n = 1, \dots, N. \quad (16)$$

The space $(\mathcal{H}_N)_1$ is N -dimensional:

$$(\mathcal{H}_N)_1 = \bigoplus_{1 \leq n \leq N} \mathbb{C}|n\rangle, \quad |n\rangle = \sigma_n^- |\Omega_+\rangle = |\uparrow\dots\uparrow \underset{n}{\downarrow} \uparrow\dots\uparrow\rangle. \quad (17)$$

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$$H_{\text{XXZ}}|z\rangle = \left(-\frac{N\Delta}{2} + \epsilon(z) \right) |z\rangle, \quad \epsilon(z) = 2\Delta - z - z^{-1}. \quad (22)$$

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Consider the case $k = 2$. Let us search for an eigenstate in the form of a combination of wave solutions:

$$|z_1, z_2\rangle = \sum_{n_1 < n_2} (A_{12} z_1^{n_1} z_2^{n_2} + A_{21} z_2^{n_1} z_1^{n_2}) |n_1, n_2\rangle. \quad (23)$$

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$$\frac{A_{21}}{A_{12}} = S(z_1, z_2) \equiv -\frac{1 + z_1 z_2 - 2\Delta z_2}{1 + z_1 z_2 - 2\Delta z_1}. \quad (24)$$

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These two equations determine the allowed values of the pairs (z_1, z_2) . Thus, though the total energy looks additive, the sets of allowed quasimomenta of excitations are different in the cases of one-particle and of two-particle states.

Consider general k . The **Bethe Ansatz** is

$$|z_1, \dots, z_k\rangle = \sum_{n_1 < \dots < n_k} \sum_{\sigma \in S_k} A_{\sigma_1 \dots \sigma_k} \prod_{j=1}^k z_{\sigma_j}^{n_j} |n_1, \dots, n_k\rangle.$$

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$$H_{\text{XXZ}} |\Psi_k(z_1, \dots, z_k)\rangle = \left(-\frac{N\Delta}{2} + \sum_{i=1}^k \epsilon(z_i) \right) |\Psi_k(z_1, \dots, z_k)\rangle, \quad (28)$$